

Mathematics

ON SOME SINGULAR INTEGRAL EQUATIONS ON THE SEMI-AXIS

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In this paper the non-characteristic singular integral equations on the semi-axis are discussed. The solution of these equations is reduced to the solution of one-dimensional pseudo-differential equations. Some examples of singular equations for which explicit solutions exist are provided.

**Keywords:** singular equation, pseudo-differential equation, matrix coupling.

**1<sup>0</sup>.** Let  $\overline{\mathbb{R}}$  be a two-point compactification of  $\mathbb{R} = (-\infty, +\infty)$  and  $C(\overline{\mathbb{R}})$  – the Banach algebra of all complex valued continuous function on  $\overline{\mathbb{R}}$ . What follows,  $\rho(z) = z^\beta$  ( $-1 < \beta < 1$ ) is understood as the branch of this function that is analytical in  $\mathbb{C} \setminus (-\infty, 0)$  and assumes positive values on the positive semi-axis  $\mathbb{R}_+ = (0, +\infty)$ . Let  $\Sigma_\beta$  be a subalgebra of all linear bounded operators acting from and to the weighted space  $L_2(\mathbb{R}_+, \rho)$ , generated by operators  $I$  and  $S_{\mathbb{R}_+}$ , where  $S_{\mathbb{R}_+}$  is a singular integral operator acting along  $\mathbb{R}_+$ :

$$(S_{\mathbb{R}_+} y)(x) = \frac{1}{\pi i} \int_0^{+\infty} \frac{y(\tau)}{\tau - x} d\tau.$$

Here the integral is perceived in terms of the main value. The notation  $F$  means the following Fourier transform:

$$(Fy)(\xi) = \hat{y}(\xi) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ix\xi} y(x) dx \quad (y \in L_2(\mathbb{R}), \xi \in \mathbb{R}),$$

and  $\gamma$  is the operator continuously reflecting  $L_2(\mathbb{R}_+, \rho)$  in  $L_2(\mathbb{R})$  according to formula  $(\gamma y)(x) = e^{sx} y(e^{\alpha x})$ , where  $\alpha > 0$ ,  $\sigma = \alpha(\beta + 1)/2$  and  $s = \sigma + i\xi$ . It is not difficult to ascertain that the operator  $\gamma^{-1}F$  is the Mellin transform.

Below the multiplicative operator by function  $u$  is denoted as  $A_u$  (i.e.  $A_u y = uy$ ). For any function  $a \in C(\overline{\mathbb{R}})$  the operator  $K_a = \gamma^{-1}F A_u F^{-1}\gamma$  belongs to algebra  $\Sigma_\beta$  (see, e.g., [1, 2]).

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As is well known, the characteristic equation  $(A_a + S_{\mathbb{R}_+} A_b)y = f$  ( $y, f \in L_2(\mathbb{R}_+, \rho)$ ) admits the explicit solution [3]. The present paper is devoted to the investigation of “difficult” singular integral operators of a type  $\tilde{K} = \sum_{m=1}^N K_{\varphi_m} A_{\psi_m}$  ( $\varphi_m \in C(\overline{\mathbb{R}}), \psi_m \in L_\infty(\mathbb{R}_+)$ ), acting from  $L_2(\mathbb{R}_+, \rho)$  to itself.

Let  $\mathbb{H}_r(\mathbb{R})$  ( $r \geq 0$ ) be Sobolev–Slobodetski spaces of the generalized functions  $u$ , the Fourier transform  $\hat{u}$  of which belongs to  $L_2(\mathbb{R}, (1+|x|)^r)$  space. Following [4], the class of locally integrable functions  $A$  on  $\mathbb{R}$ , complying to condition  $|A(\xi)| \leq c(1+|\xi|)^r$ , is denoted as  $\mathbb{S}_r^0$ .

Let  $\psi_0(x) = (1 + |\alpha^{-1} \ln x|)^r$  ( $x \in \mathbb{R}_+$ ) and functions  $\psi_i \in L_\infty(\mathbb{R}_+)$  ( $i = 1, \dots, N$ ). Now determine the functions  $A_0 = \psi_0 \circ l$ ,  $A_m = (\psi_m \circ l) A_0$  ( $m = 1, \dots, N$ ) on  $\mathbb{R}$ , where  $l(x) = e^{\alpha x}$  ( $x \in \mathbb{R}$ ). It follows from here that  $\psi_0 = A_0 \circ h$  and  $\psi_m = (A_m \circ h) \psi_0^{-1}$  ( $m = 1, \dots, N$ ), where  $h(x) = \alpha^{-1} \ln x$  when  $x \in \mathbb{R}_+$ . It is obvious that  $A_m \in \mathbb{S}_r^0$  for all  $m = 0, \dots, N$ .

Let us consider a operator  $A(x, D)$  with a symbol  $A(x, t) = \sum_{m=1}^N \varphi_m(x) A_m(t)$ :

$$A(x, D)u = \int_{-\infty}^{\infty} e^{-ixt} A(x, t) \hat{u}(t) dt.$$

Since  $A(x, D)u = \sum_{m=1}^N A_{\varphi_m} A_m(D)u$ , then for  $r \geq 0$  the reflection  $A(x, D)$  may be prolonged till the continuous reflection from  $\mathbb{H}_r(\mathbb{R})$  to  $L_2(\mathbb{R})$  (see [4]).

Below the solution of “difficult” equation  $\tilde{K}z = g$  is reduced to a pseudo-differential equation  $A(x, D)y = f$  for some  $f$ . The examples of “difficult” singular integral equations investigated based on this relationship are provided.

**2<sup>0</sup>.** Let  $X_i$  ( $i = 1, \dots, 4$ ) be linear spaces on the field of complex numbers,  $\omega_1: X_1 \rightarrow X_2$ ,  $\omega_2: X_2 \rightarrow X_1$ ,  $\omega_3: X_1 \rightarrow X_3$ ,  $\omega_4: X_4 \rightarrow X_1$  are linear reflections and  $\omega_4$  is an invertible reflection. Let us consider the linear reflections

$$\begin{aligned} T: X_2 &\rightarrow X_2 \oplus X_3, & K: X_4 &\rightarrow X_1 \oplus X_3, \\ A_{12}: X_1 \oplus X_3 &\rightarrow X_2 \oplus X_3, & A_{21}: X_2 &\rightarrow X_4, \\ A_{22}: X_1 \oplus X_3 &\rightarrow X_4, & B_{11}: X_2 \oplus X_3 &\rightarrow X_2, \\ B_{12}: X_4 &\rightarrow X_2, & B_{21}: X_2 \oplus X_3 &\rightarrow X_1 \oplus X_3, \end{aligned}$$

determined by equations

$$T = \begin{bmatrix} I_{X_2} + \omega_1 \omega_2 \\ \omega_3 \omega_2 \end{bmatrix}, \quad A_{12} = \begin{bmatrix} -\omega_1 & 0 \\ -\omega_3 & I_{X_3} \end{bmatrix},$$

$$A_{21} = [-\omega_4^{-1}\omega_2], \quad A_{22} = [\omega_4^{-1} \quad 0], \quad B_{11} = [I_{X_2} \quad 0], \quad B_{12} = [\omega_1\omega_4],$$

$$B_{21} = \begin{bmatrix} \omega_2 & 0 \\ 0 & I_{X_3} \end{bmatrix}, \quad K = \begin{bmatrix} \omega_4 + \omega_2\omega_1\omega_4 \\ \omega_3\omega_4 \end{bmatrix}.$$

By means of direct calculations it is easy to make sure of the validity of equality

$$\begin{bmatrix} T & A_{12} \\ A_{21} & A_{22} \end{bmatrix}^{-1} = \begin{bmatrix} B_{11} & B_{12} \\ B_{21} & K \end{bmatrix}.$$

Based on this equality and results of [5] the following statement is derived.

*Assignment 1.* If  $K^{(-1)}$  is pseudo-inverse of  $K$ , then the reflection

$$T^{(-1)} = B_{11} - B_{12}K^{(-1)}B_{21}$$

is pseudo-inverse of  $T$ . Similarly, if  $T^{(-1)}$  is pseudo-inverse of  $T$ , then

$$K^{(-1)} = A_{22} - A_{21}T^{(-1)}A_{12}$$

is pseudo-inverse of  $K$ . Spaces  $\text{Ker}T$  and  $\text{Ker}K$  are isomorphic. Besides, the following equalities are valid:

$$\text{Ker}T = B_{12}\text{Ker}K, \quad \text{Ker}K = A_{21}\text{Ker}T,$$

$$\text{Im}T = B_{21}^{-1}\text{Im}K, \quad \text{Im}K = A_{12}^{-1}\text{Im}T.$$

**3<sup>o</sup>.** Let  $\mathbb{M}_r(\mathbb{R})$  be some direct addition to the linear space  $\mathbb{H}_r(\mathbb{R})$  in  $L_2(\mathbb{R})$  space, and

$$\pi_1 : L_2(\mathbb{R}) \rightarrow \mathbb{H}_r(\mathbb{R}), \quad \pi_2 : L_2(\mathbb{R}) \rightarrow \mathbb{M}_r(\mathbb{R})$$

are projection operators that are related by the ratio

$$\pi_1 y + \pi_2 y = y, \quad y \in L_2(\mathbb{R}).$$

Let us determine a space  $W = \{f : \psi_0^{-1}f \in L_2(\mathbb{R}_+, \rho)\}$ . Then, apply the Assignment 1 for spaces

$$X_1 = L_2(\mathbb{R}_+, \rho), \quad X_2 = \mathbb{H}_r(\mathbb{R}), \quad X_3 = \mathbb{M}_r(\mathbb{R}), \quad X_4 = W$$

and operators

$$\omega_1 = \pi_1 F^{-1}\gamma, \quad \omega_2 = \sum_{m=1}^N K_{\varphi_m} \gamma^{-1} F A_m(D) - \gamma^{-1} F, \quad \omega_3 = \pi_2 F^{-1}\gamma, \quad \omega_4 = A_{\psi_0^{-1}}.$$

By using the equality  $F^{-1}\gamma K_{\varphi_m} \gamma^{-1} F = A_{\varphi_m}$  ( $m=1, \dots, N$ ) it is easy to ascertain that the operator  $T$  acting from  $\mathbb{H}_r(\mathbb{R})$  to  $L_2(\mathbb{R})$  coincides with  $A(x, D)$ .

**Lemma 1.** Let

$$f = (\tilde{f}, 0) \in L_2(\mathbb{R}_+, \rho) \oplus \mathbb{M}_r(\mathbb{R}).$$

The function  $z$  is the solution of equation  $Kz = f$ , iff  $z \in L_2(\mathbb{R}_+, \rho)$  and  $\tilde{K}z = \tilde{f}$ .

*Proof.* Let  $z$  be the solution of equation  $Kz = f$ , that is

$$\omega_4 z + \omega_2 \omega_1 \omega_4 z = \tilde{f}, \quad \omega_3 \omega_4 z = 0. \quad (1)$$

First prove that  $z \in L_2(\mathbb{R}_+, \rho)$ . Really, from the second equation of (1) it follows that  $g = F^{-1}\gamma A_{\psi_0^{-1}} z \in \mathbb{H}_r(\mathbb{R})$ . Consequently,  $\gamma z = F A_0(D)g \in L_2(\mathbb{R})$ ,

that is  $z \in L_2(\mathbb{R}_+, \rho)$ . Note now that from (1) it also follows that  $\pi_1 F^{-1} \gamma A_{\psi_0^{-1}} z = F^{-1} \gamma A_{\psi_0^{-1}} z$ . From the first equation of (1) it follows that  $\tilde{K}z = \tilde{f}$ .

Let  $z \in L_2(\mathbb{R}_+, \rho)$  and  $\tilde{K}z = \tilde{f}$ . From equality  $F^{-1} \gamma A_{\psi_0^{-1}} z = A_0^{-1}(D) F^{-1} \gamma z$  it follows that

$$\omega_3 \omega_4 z = 0.$$

Since  $\tilde{K}z = \tilde{f}$ , we have

$$\begin{aligned} \left( \sum_{m=1}^N K_{\varphi_m} A_{\psi_m} \right) z &= \left( \sum_{m=1}^N K_{\varphi_m} A_{A_m \circ h} A_{\psi_0^{-1}} \right) z = \left( A_{\psi_0^{-1}} + \left( \sum_{m=1}^N K_{\varphi_m} A_{A_m \circ h} - I \right) A_{\psi_0^{-1}} \right) z = \\ &= \left( A_{\psi_0^{-1}} + \left( \sum_{m=1}^N K_{\varphi_m} \gamma^{-1} F A_m(D) - \gamma^{-1} F \right) F^{-1} \gamma A_{\psi_0^{-1}} \right) z = \\ &= \left( A_{\psi_0^{-1}} + \left( \sum_{m=1}^N K_{\varphi_m} \gamma^{-1} F A_m(D) - \gamma^{-1} F \right) \pi_1 F^{-1} \gamma A_{\psi_0^{-1}} \right) z = \\ &= (\omega_4 + \omega_2 \omega_1 \omega_4) z = \tilde{f}. \end{aligned}$$

Thus, the equalities are correct, that is  $Kz = f$ . The proof is completed.

**Theorem 1.** To ensure that function  $z \in L_2(\mathbb{R}_+, \rho)$  satisfies the equation

$$\sum_{m=1}^N K_{\varphi_m} A_{\psi_m} z = g \quad (g \in L_2(\mathbb{R}_+, \rho)), \quad (2)$$

it is necessary and sufficient that function  $y = F^{-1} \gamma A_{\psi_0^{-1}} z$  satisfy the equation

$$A(x, D)y = F^{-1} \gamma g \quad (F^{-1} \gamma g \in L_2(\mathbb{R})). \quad (3)$$

*Proof.* Let the function  $z \in L_2(\mathbb{R}_+, \rho)$  be a solution of equation (2). Acting by operator  $F^{-1} \gamma$  on both the parts of the equation (2) (it is feasible, since the functions  $\psi_m$  ( $m=1, \dots, N$ ) are bounded), we obtain

$$\sum_{m=1}^N F^{-1} \gamma K_{\varphi_m} A_{A_m \circ h} A_{\psi_0^{-1}} z = F^{-1} \gamma g.$$

Taking into account the equalities  $A_{A_m \circ h} = \gamma^{-1} F A_m(D) F^{-1} \gamma$  and  $F^{-1} \gamma K_{\varphi_m} \gamma^{-1} F = A_{\varphi_m}$  ( $m=1, \dots, N$ ), we find

$$\sum_{m=1}^N A_{\varphi_m} A_m(D) y = F^{-1} \gamma g.$$

Earlier, it has been shown that  $y = F^{-1} \gamma A_{\psi_0^{-1}} z \in \mathbb{H}_r(\mathbb{R})$  (see the proof of Lemma 1). Vice versa, now let the function  $y \in \mathbb{H}_r(\mathbb{R})$  be the solution of equation (3). Then it is true that  $z = A_{\psi_0} \gamma^{-1} F y \in L_2(\mathbb{R}_+, \rho)$ , the proof of which is provided in Lemma 1.

At the insertion of  $y = F^{-1} \gamma A_{\psi_0^{-1}} z$  into (3), we have

$$\sum_{m=1}^N A_{\varphi_m} A_m(D) F^{-1} \gamma A_{\psi_0^{-1}} z = F^{-1} \gamma g.$$

Taking into account the equality  $A_{\varphi_m} = F^{-1}\gamma K_{\varphi_m}\gamma^{-1}F$  ( $m=1, \dots, N$ ), we obtain

$$\sum_{m=1}^N F^{-1}\gamma K_{\varphi_m}\gamma^{-1}FA_m(D)F^{-1}\gamma A_{\psi_0^{-1}}z = F^{-1}\gamma g. \quad (4)$$

Acting by operator  $\gamma^{-1}F$  on both the parts of equality (4), we find

$$\sum_{m=1}^N K_{\varphi_m}\gamma^{-1}FA_m(D)F^{-1}\gamma A_{\psi_0^{-1}}z = g.$$

By using equality  $\gamma^{-1}FA_m(D)F^{-1}\gamma = A_{A_m \circ h}$  ( $m=1, \dots, N$ ), we find that

$\sum_{m=1}^N K_{\varphi_m} A_{\psi_0^{-1}}z = g$ , i.e. the function  $z$  is the solution of equation (2).

The Theorem is proved.

**4<sup>0</sup>.** In case of  $N=2$ ,  $\varphi_1(x)=1$  and  $\varphi_2(x) = \left[ \operatorname{ch} \frac{\pi}{\alpha} \left( x - \xi + i \frac{\alpha\beta}{2} \right) \right]^{-2}$  we

obtain  $K_{\varphi_1} = I$  and  $(K_{\varphi_2}y)(x) = \frac{1}{\pi^2} \int_0^{+\infty} y(\xi) \frac{\ln \xi - \ln x}{\xi - x} d\xi$ .

Now, let us highlight two examples.

*Example 1.* Let  $A_1(\xi) = -\xi^2 - \frac{\pi^2}{\alpha^2}$ ,  $A_2(\xi) = 2\frac{\pi^2}{\alpha^2}$  and  $r=2$ , then we have

$$A(x, D)y(x) = y''(x) + \frac{\pi^2}{\alpha^2} \left( 2 \left[ \operatorname{ch} \frac{\pi}{\alpha} \left( x - \xi + i \frac{\alpha\beta}{2} \right) \right]^{-2} - 1 \right) y(x),$$

$$\widetilde{K}z(x) = \frac{-(\alpha^{-1} \ln x)^2 - \pi^2/\alpha^2}{(1 + |\alpha^{-1} \ln x|)^2} z(x) + \frac{2}{\alpha^2} \int_0^{+\infty} \frac{z(\xi)}{(1 + |\alpha^{-1} \ln \xi|)^2} \cdot \frac{\ln \xi - \ln x}{\xi - x} d\xi.$$

Taking into account that functions

$$y_1(x) = e^{\frac{\pi}{\alpha}x} \left[ 1 + e^{\frac{2\pi}{\alpha} \left( x - \xi + i \frac{\alpha\beta}{2} \right)} \right]^{-1}, \quad y_2(x) = e^{-\frac{\pi}{\alpha}x} \left[ 1 + e^{-\frac{2\pi}{\alpha} \left( x - \xi + i \frac{\alpha\beta}{2} \right)} \right]^{-1}$$

are solutions of equation  $A(x, D)y = 0$  from class  $\mathbb{H}_2(\mathbb{R})$ , we obtain according to the Assignment 1 that functions  $z_k(x) = -\omega_4^{-1} \omega_2 y_k(x)$  ( $k=1, 2$ ) are the basis of  $\operatorname{Ker} \widetilde{K}$ .

*Example 2.* Let  $A_1(\xi) = -a_1 i \xi + a_2$ ,  $A_2(\xi) = b_1 i \xi - b_2$  and  $r=1$ , where  $a_1, a_2, b_1, b_2 \in \mathbb{C}$ , then we find

$$A(x, D)y(x) = \left( a_1 - b_1 \left[ \operatorname{ch} \frac{\pi}{\alpha} \left( x - \xi + i \frac{\alpha\beta}{2} \right) \right]^{-2} \right) y'(x) + \left( a_2 - b_2 \left[ \operatorname{ch} \frac{\pi}{\alpha} \left( x - \xi + i \frac{\alpha\beta}{2} \right) \right]^{-2} \right) y(x),$$

$$\widetilde{K}z(x) = \frac{-a_1 i \alpha^{-1} \ln x + a_2}{1 + |\alpha^{-1} \ln x|} z(x) - \frac{1}{\pi^2} \int_0^{+\infty} \frac{-b_1 i \alpha^{-1} \ln \xi + b_2}{1 + |\alpha^{-1} \ln \xi|} \cdot \frac{\ln \xi - \ln x}{\xi - x} z(\xi) d\xi.$$

Taking into account that the only solution of equation  $A(x, D)y(x) = 0$  does not belong to class  $L_2(\mathbb{R})$ , we obtain that equation  $\tilde{K}z(x) = 0$  does not have any solutions from  $L_2(\mathbb{R}_+, \rho)$ .

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Վ. Վ. Մինոնյան

Կիսասառանգքի վրա որոշ սինգուլյար ինտեգրալ հավասարումների մասին

Աշխատանքում դիտարկվում են սինգուլյար ինտեգրալ հավասարումներ կիսասառանգքի վրա: Այս հավասարումների լուծումը հանգեցվում է միաչափ պսևդոդիֆերենցիալ հավասարումների լուծմանը: Բերված են բացահայտ տեսքով լուծելի սինգուլյար հավասարումների օրինակներ:

В. В. Симонян

О некоторых сингулярных интегральных уравнениях на полуоси

В работе рассматриваются нехарактеристические сингулярные интегральные уравнения на полуоси. Решение этих уравнений сводится к решению одномерных псевдодифференциальных уравнений. Приведены примеры сингулярных уравнений, разрешимых в явном виде.